

Fredrickson-Andersen 2 spin facilitated model

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joint with **Ivailo Hartarsky** (TU Wien) and **Fabio Martinelli** (Univ. Roma Tre)

The model

An interacting particle system on $\{0, 1\}^{\mathbb{Z}^d} = \{\circ, \bullet\}^{\mathbb{Z}^d}$, $d \geq 2$.

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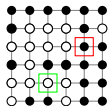
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- the refresh occurs **iff** the site has **at least 2 empty nearest neighbours** = *iff the kinetic constraint is satisfied*



FA-2f: properties

- Reversible w.r.t. Bernoulli(1-q) product measure, μ_q

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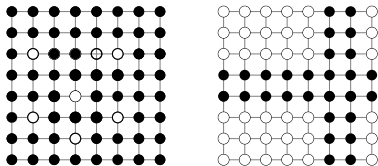
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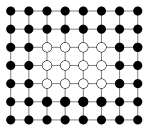


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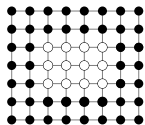
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Math. motivation #1: several IPS tools fail → new tools needed!

Motivations from physics

- FA2f was introduced in the '80's to model/understand the **liquid/glass transition**, a major open problem in physics
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→ **no clear-cut answers / contradicting conjectures**
- **Math. motivation #2**: put physicists works on firmer ground / settle controversies

Does $\eta \sim \mu_q$ contain a blocked cluster?

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How does the first time the origin is empty scale as $q \downarrow 0$?

$$\tau^{\text{BP}} = \exp\left(\lambda_d q^{-1/(d-1)}(1 + o(1))\right) \quad [\text{Aizenman - Lebowitz '88}]$$

- $\lambda_2 = \pi^2/18$ [Holroyd '08]
- $\lambda_d = \dots \forall d > 2$ [Balogh Bollobas Duminil-Copin Morris '12]

Back to FA2f: our results

Theorem [Hartarsky, Martinelli, C.T. '20]

As $q \downarrow 0$, w.h.p. for the stationary FA-2f model on \mathbb{Z}^d it holds

$$\tau = \exp \left(\frac{d \times \lambda_d}{q^{1/(d-1)}} (1 + o(1)) \right), \quad d \geq 2$$

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Remark

- not a corollary of BP result: very different mechanism!
- we settle contrasting conjectures in physics literature

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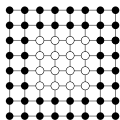
- Step 1* make a good guess for the **optimal relaxation mechanism**
- Step 2* develop a toolbox (**Poincaré inequalities + renormalisation**) to translate heuristics into rigorous bounds $\implies \tau \leq \dots$
- Step 3* identify a **bottleneck**, i.e. an unlikely configuration set that has to be visited before emptying the origin $\implies \tau \geq \dots$

Heuristics ($q \downarrow 0$): optimal relaxation mechanism

- Relaxation is driven by the motion of rare large patches of empty sites, the droplets
- droplet density $\rho_D := \exp\left(-\frac{d \times \lambda_d}{q^{1/d-1}}(1 + o(1))\right)$
droplet length $L_D := 1/q^\alpha, \alpha > 2$
- Droplets can move in any direction

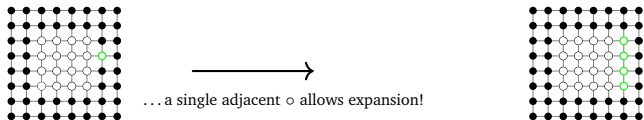
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- **Droplets can move in any direction ... isn't this contradictory with "finite empty regions cannot expand"?!**



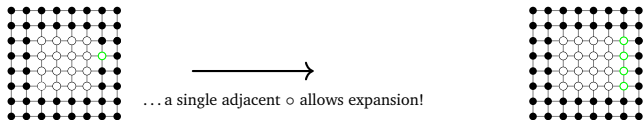
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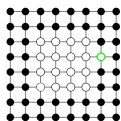
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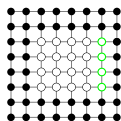


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→
... a single adjacent \circ allows expansion!



- Droplet motion requires **few additional empty sites** \rightarrow
this **good environment is very likely** since $L_D \gg |\log q|/q$

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$\tau^{\text{BP}} \sim$ size of minimal region to be unblocked before the origin
= distance from the origin to the nearest droplet

$$\rightarrow \tau^{\text{BP}} \sim \rho_D^{-1/d} \sim \tau^{1/d}$$

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- the supremum is over non constant functions;
- $\mathcal{D}_{\text{FA-2f}}(f)$ is the "energy" $= \sum_\eta \mu(\eta) \sum_x c_x(\eta) (f(\eta^x) - f(\eta))^2$
- η^x = configuration flipped at x ; $c_x(\eta)$ = rate for $\eta \rightarrow \eta^x$

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- Renormalisation on droplet scale \rightarrow CBSEP dynamics for the droplets and FA2f dynamics inside the droplets

$$T_{\text{FA-2f}}^{\text{rel}} \leq T_{\text{CBSEP}}^{\text{rel}} \times T_{\text{FA-2f}}^{\text{rel}}(L_D \mid \text{droplet})$$

- Prove Poincaré inequalities for FA2f in a droplet and CBSEP

$$T_{\text{CBSEP}}^{\text{rel}} \leq \rho_D^{-1}, \quad T_{\text{FA-2f}}^{\text{rel}}(L_D \mid \text{droplet}) \leq e^{\frac{|\log q|^3}{\sqrt{q}}} \ll T_{\text{CBSEP}}^{\text{rel}}$$

Why bothering with the exact constants?

- sharp constant is extremely hard to grasp numerically:
 - subtle corrections to the dominant behavior, ex. $d = 2$:
 $\tau^{\text{BP}} = \exp\left(\frac{\pi^2}{18q}(1 - c\sqrt{q})\right)$ (Hartarsky, Morris)
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- a deeper understanding of the cooperative relaxation
- the mathematical tools we build are **very flexible**
→ we adapt them to get **universality results for KCM in $d = 2$**

Open problems

- **Conjecture #1:** Start from $\mu_{q'}$, with $q' \neq q$ and $q, q' > 0$ and call μ^t the evolved measure at time t , $\mu^t = \mu_{q'} P_t$. It holds

$$\lim_{t \rightarrow \infty} \mu^t = \mu_q.$$

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More generally: we lack robust tools to tackle the out of equilibrium regime of KCM !

Thanks for your attention!

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- Super-good (SG) rectangles:
 - a rectangle of class 0 is SG if it is empty;
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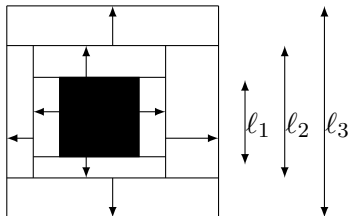
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Droplets are defined as $\ell_N \times \ell_N$ SG rectangles

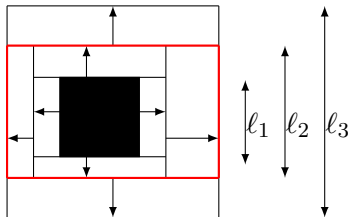
How do droplets look like?

Ex. of a SG rectangle of class 6. Here arrows indicate traversability and the black square is completely empty.



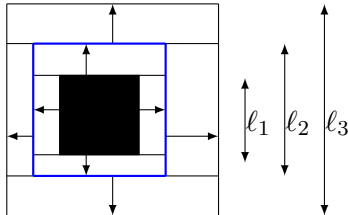
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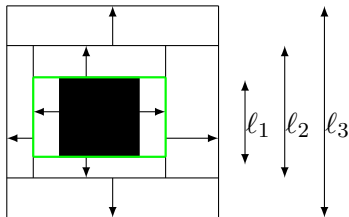
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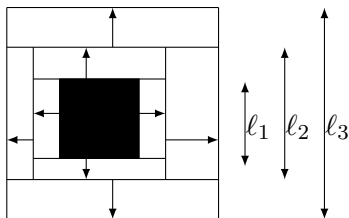
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Droplets = $\ell_N \times \ell_N$ squares that are SG with

$N = 8 \lfloor \log q \rfloor / \sqrt{q}$, so that $\ell_N = q^{-17/2+o(1)}$

NB any other big enough power would work the same, no special meaning of $17/2$...

FA- j f model

- for j n-BP for all $d \geq j \geq 2$, w.h.p. it holds

$$\tau_0^{\text{BP}} \sim \exp^{j-1} \left(\frac{\tilde{\lambda}_{d,j}}{q^{1/(d-j+1)}} \right)$$

\exp^k = exponential iterated k times (Balogh, Bollobas, Duminil-Copin, Morris '12)

Same scaling for τ_0 (Hartarsky, Martinelli, C.T. in progress)

- $j = 1$: $\tau_0^{\text{BP}} = 1/q^{1/d}$, $\tau_0 = 1/q^{\nu(d)}$,
 $\nu(1) = 3$, $\nu(d) = 2$ $d \geq 2$ (log corrections in $d = 2$)
(Cancrini, Roberto, Martinelli, C.T. '08 + Shapira '20)
- $d < j$: $\tau_0 = \tau_0^{\text{BP}} = \infty$ w.h.p. for $q \rightarrow 0$

What happens if we change constraint?

Universality results for $d = 2$

- 1 **Supercritical unrooted:** $\tau(q) = q^{-\Theta(1)}$
- 2 **Supercritical rooted:**
 $\tau(q) = q^{-\Theta(1)|\log q|}$
- 3 **Finitely critical:** $\tau(q) = \exp\left(\frac{\Theta(1)(\log q)^{\Theta(1)}}{q^\nu}\right)$
- 4 **Infinitely critical:**
 $\tau(q) = \exp(q^{-2\nu}(\log q)^c)$
- 5 **Subcritical:**
 $\exists q_c > 0$, s.t. for $q < q_c$ it holds $\tau(q) = \infty$

[I.Hartarsky, L.Marêché, F.Martinelli, R. Morris, C.T.]

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- 2 **Supercritical rooted:**
 $\tau(q) = q^{-\Theta(1)|\log q|} \gg \tau^{\text{BP}} = q^{-\Theta(1)}$ (East)
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(Duarte)
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[I.Hartarsky, L.Marêché, F.Martinelli, R. Morris, C.T.]

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logarithmic energy barrier: droplet are at distance $\ell \sim \tau^{\text{BP}}$ from the origin and must create $\sim \log \ell$ droplets to reach it [Marêché, Martinelli, C.T. '20]

The g -CBSEP chain

- $G = (V, E)$: finite connected graph
- (\mathcal{S}, π) : finite probability space
- $\mathcal{S} = \mathcal{S}_0 \sqcup \mathcal{S}_1$ and $\rho = \pi(\mathcal{S}_1)$
- given $\sigma \in \mathcal{S}^V$, $x \in V$ is *occupied* iff $\sigma_x \in \mathcal{S}_1$
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Theorem [Hartarsky, Martinelli, C.T. '20]

As $\rho \downarrow 0$, $T_{\text{rel}}^{\text{g-CBSEP}} \leq O(\rho^{-1} \log(1/\rho))$